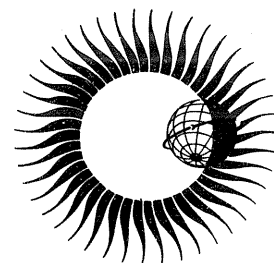


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DESCRIPTION AND CATALOG
OF IONOSPHERIC F-REGION DATA,
JICAMARCA RADAR OBSERVATORY
(NOVEMBER 1966 - APRIL 1969)

April 1976



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REPORT UAG - 53

DESCRIPTION AND CATALOG OF IONOSPHERIC F-REGION DATA, JICAMARCA RADAR OBSERVATORY (NOVEMBER 1966 - APRIL 1969)

by

W. L. Clark¹, J. P. McClure² and T. E. VanZandt¹

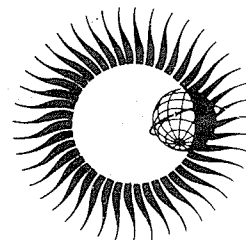
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April 1976

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(NOVEMBER 1966 - APRIL 1969)

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Introduction

This report basically describes equatorial ionospheric F-region data reduced from the Jicamarca Radar Observatory (JRO) incoherent scatter observations for particular periods. It lists in catalog form the times of the observations made during those periods. These F-region data include the electron concentration, Ne, and the electron and ion temperatures, Te and Ti. The data were inferred from the incoherent scatter observations of JRO which is located at S11.95 E283.13. JRO is the only incoherent scatter observatory close to the magnetic equator at 2°N magnetic dip.

For the period 11 November 1966 to 29 April 1969, Ne, Te and Ti data reduced from the JRO observations are now available both in digital form (on magnetic tape) and in graphical and tabular form (on microfilm) from the World Data Center A for Solar-Terrestrial Physics. Some results of similar analyses are already published, in particular the Ne data were presented in graphical form by McClure *et al.* [1970] for the period 27 May 1964 to 1 December 1966. Small amounts of data have appeared in other published papers and reports. Ionospheric drift velocity measurements reduced from the JRO incoherent scatter observations for the period July 1967 to March 1970 are found in an earlier UAG Report, Report UAG-17 [Balsley *et al.*, 1971].

Measurement Techniques and Analysis Methods

In the Faraday rotation experiment as performed at JRO, two pulses are transmitted simultaneously with opposite circular polarization. By comparing the phases of the echoes of the two pulses as a function of time delay between transmission and reception, the Faraday rotation angle as a function of height is determined. The electron concentration profile is obtained from the height derivative of the Faraday rotation angle. This technique is discussed fully by Farley [1969]. See also Evans [1969] for a general review of the incoherent scatter technique.

In addition to the measurements for Ne, the temporal autocorrelation function ρ of the radar echoes is measured at a series of time separations τ . The method by which these measurements are made is described in considerable detail by Farley *et al.* [1967] and Farley [1969] and will not be discussed here. Our concern is with the analysis of the measured autocorrelation function to deduce parameters of the ionospheric plasma.

The JRO measurements of the autocorrelation functions are usually valid only in the ionospheric F2 layer where almost all of the ions are ionized oxygen (O^+). Therefore, the ionic composition has been approximated by pure O^+ , so that the only parameters determined by the present analysis are Te and Ti. The presence of appreciable concentrations of ions other than O^+ will, of course, cause errors in the derived values of Te and Ti. In particular, near the top of the analyzed height range, especially at night during solar minimum, there may be appreciable concentrations of ionized hydrogen (H^+). For small fractional concentrations of H^+ , the fractional errors in Te and Ti are of the order of -1 and +1 times the fractional concentration, respectively. Since the fractional concentration of H^+ in the F2 layer typically increases with height with a scale height of the order of 150 km, the presence of appreciable concentrations of H^+ leads to a characteristic and obvious decrease in Te and an increase in Ti with height.

Examples of theoretical curves of $\rho(\tau)$ for pure O^+ , two values of Te and three values of Te/Ti are shown in Figure 1. Two important laws are illustrated. The first law states that the value of τ at $\rho = -0.08$ varies as $Te^{-1/2}$ with only a negligible dependence on Te/Ti. The second law states that the minimum value of ρ (ρ_{min}) in the first negative loop varies approximately linearly with Te/Ti, as illustrated in Figure 2, but is otherwise independent of Te. The following empirical formula gives Te/Ti to two decimal places:

$$Te/Ti = .2305 - 4.666 \rho_{min} + 9.23 \rho_{min}^2 \quad (1)$$

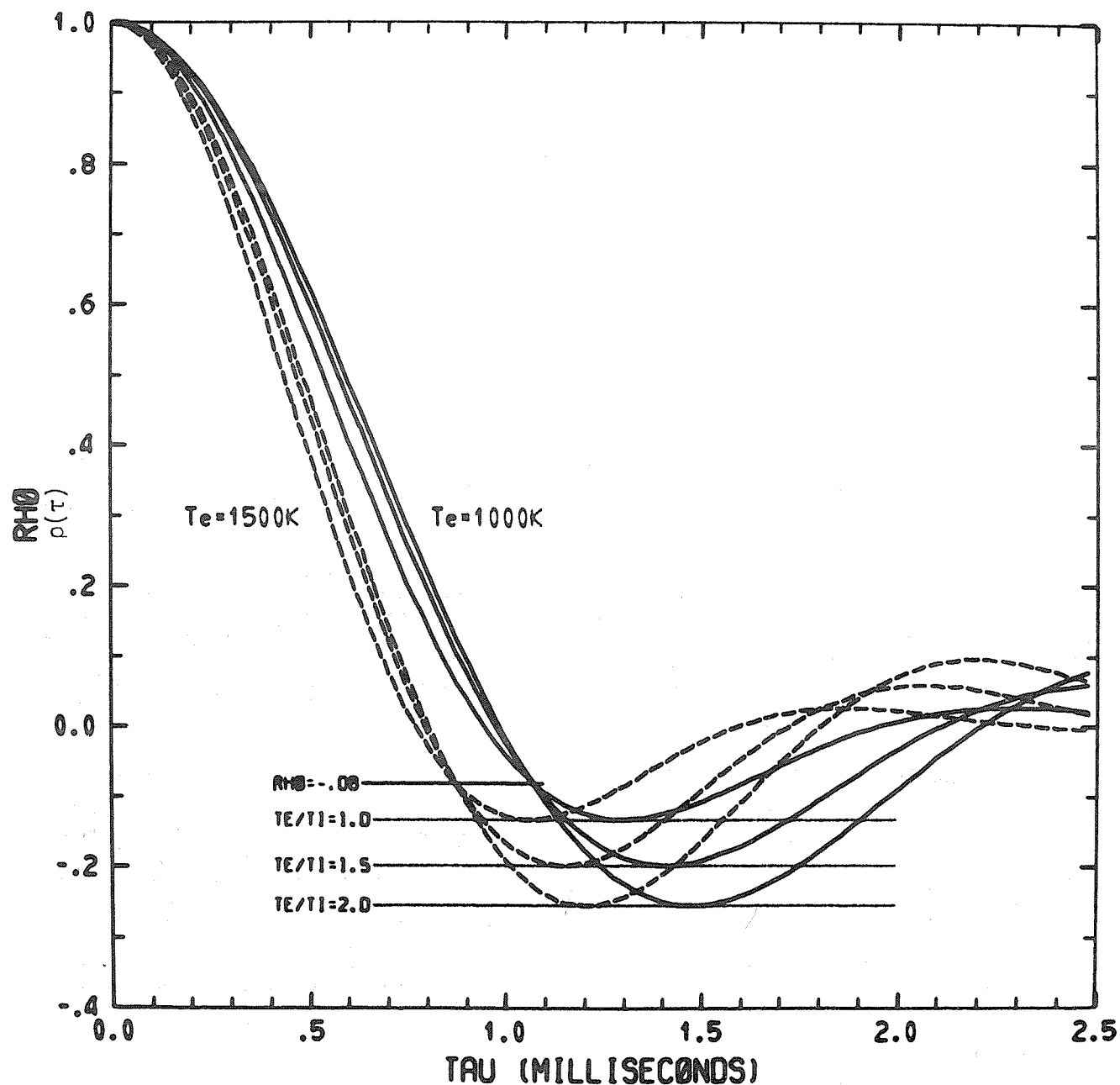


Fig. 1. Examples of the theoretical autocorrelation function $\rho(\tau)$ for $T_e = 1000K$ and $1500K$ while $T_e/T_i = 1.0, 1.5$ and 2.0 . These curves illustrate that τ at the $\rho = -0.08$ point is essentially independent of T_i and that the depth of the first negative loop depends only on T_e/T_i . These laws assume a constant ion composition, in this case pure O^+ .

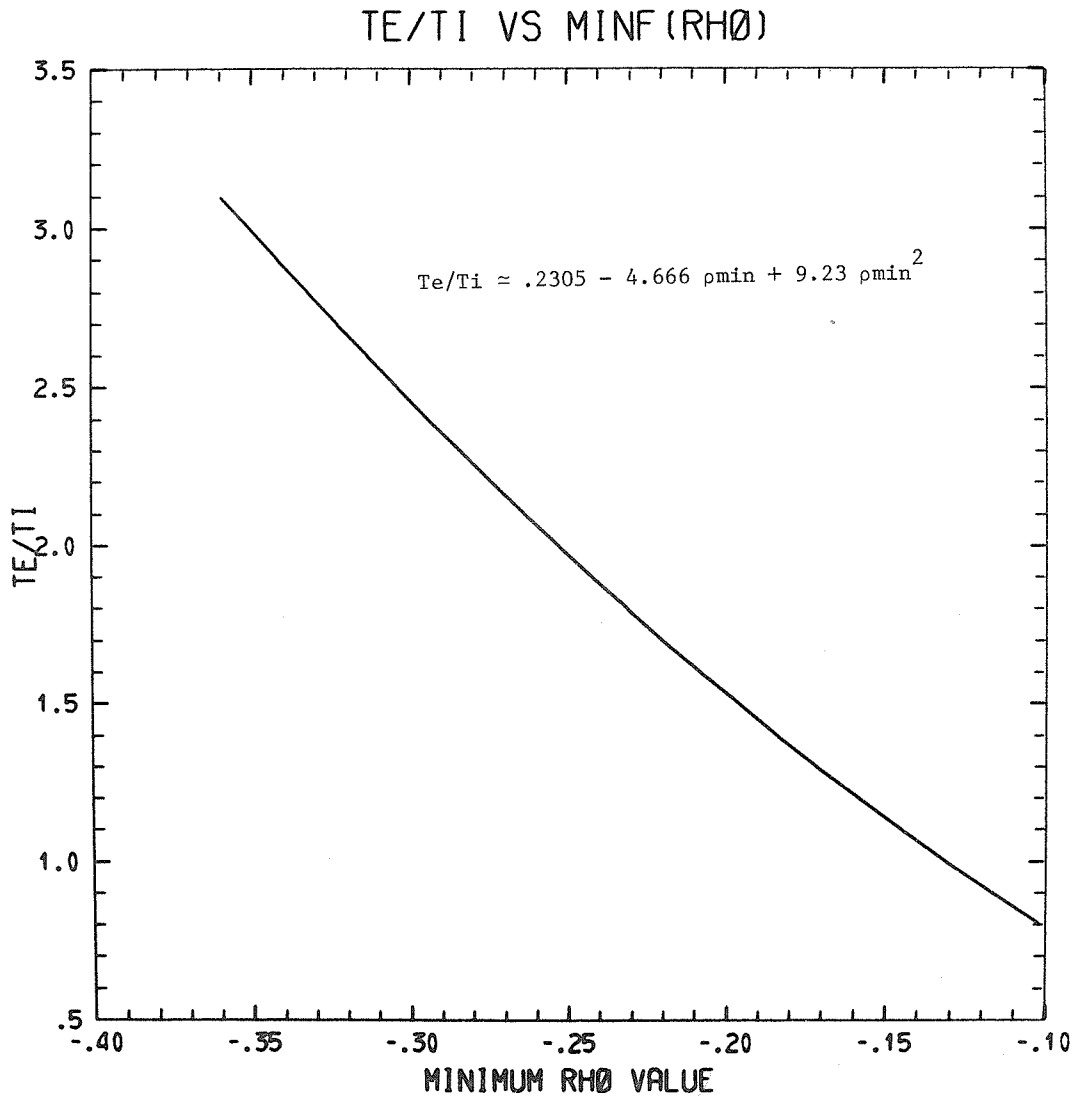


Fig. 2. Te/Ti plotted as a function of p_{min} (the minimum value of the autocorrelation function). The quadratic equation gives Te/Ti to two decimal places.

With these two laws, Te and Te/Ti may be found from the autocorrelation data by scaling the value of τ at $\rho = -0.08$ and at p_{min} . Alternatively, Te and Te/Ti can be determined by fitting the theoretical curves to the measured values of ρ . The fitting has been done in the present analysis by adjusting the values of Te and Te/Ti until the best fit in a least squares sense is obtained. This method has the advantage of using all the data and so reduces the effects of noise.

Examination of the actual autocorrelation data along with the functions fitted during the analysis can provide valuable insight into the quality of the Te and Ti values found at any particular height and time. Therefore, for each height a plot of measured ρ 's versus τ and the corresponding best-fitting function, when actually found, is included on the microfilm. A rather compact format was used to enable easy visual comparison of data from one height to the next. A sample frame is shown in Figure 3.

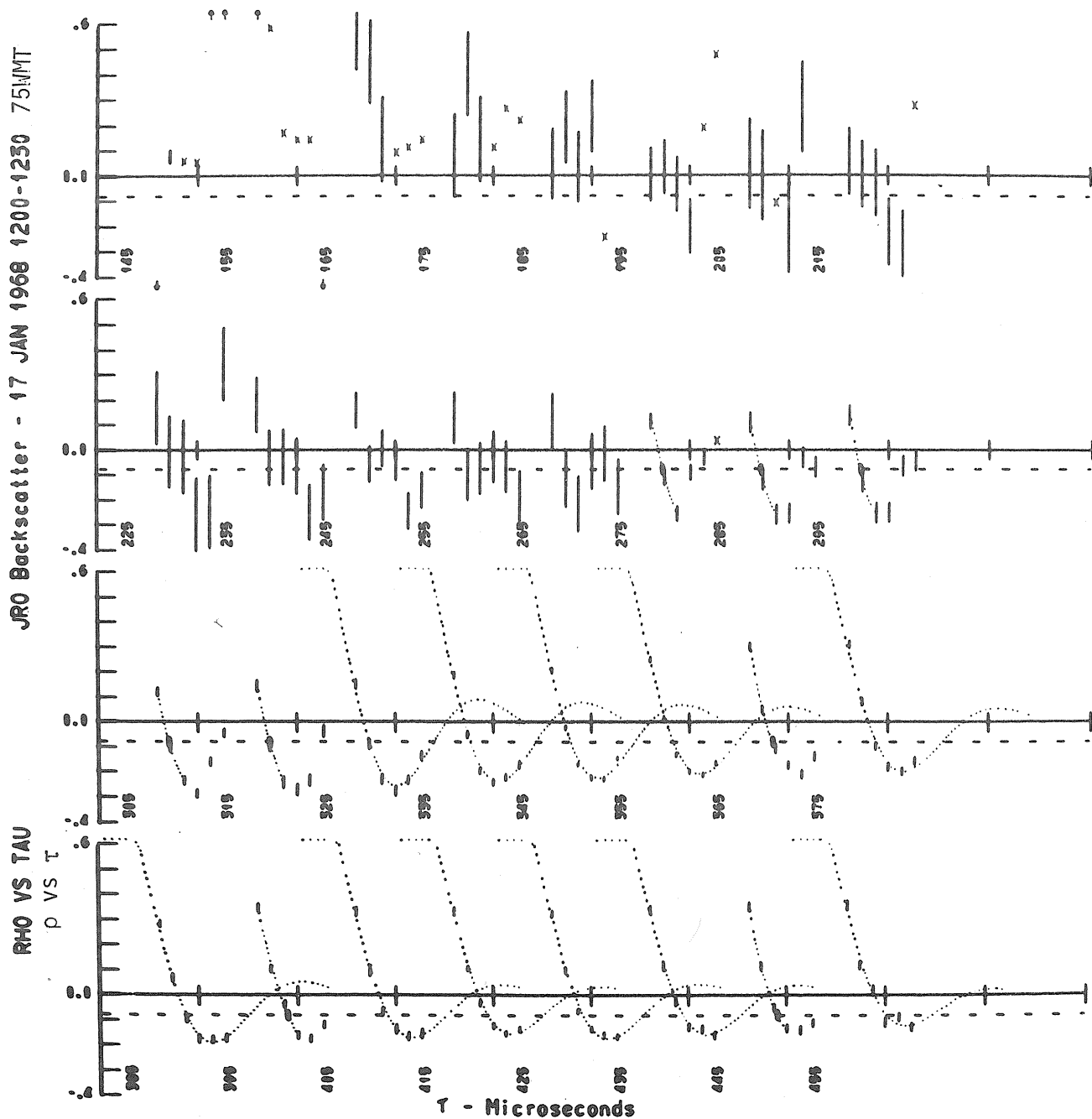


Fig. 3. A sample microfilm frame presenting the autocorrelation function found by the analysis program using the raw data and the best fitting functions. Data for 32 heights are given on each frame. This frame, for 17 Jan., 1968, 1200-1230 (75° W time), presents data for heights between 145 and 455 km in 10 kilometer steps. A much more detailed explanation of this Figure is given in the text.

Figure 3 presents data taken on 17 January 1968 during an integration from 1200 to 1230 hours (75°W time), as indicated by the title along the left margin. The Figure looks very complex because data for 32 heights are presented in overlapping form beginning with a height of 145 km and increasing in 10 km increments to 455 km. This complexity can be avoided by examining one feature at a time. Consider the bottom row of the Figure. The ordinate is the ρ axis, ranging from -0.4 to +0.6, and the abscissa is the τ axis, with a tick every 1000 μ s. Each tick mark also denotes $\tau = 0$ for the height printed sideways, below and to the right of the tick. Each measured value of ρ is plotted as a vertical bar centered on the measured value and extending $\pm \sigma$ (\pm one standard deviation). The sinuous, dotted curves in the plots for 385, 405, 415, 425, 435, and 455 km are the theoretical autocorrelation functions that best fit the measured values of ρ . The flattened part of each curve from 0 to 300 μ s is a result of an imposed plot limit just above $\rho = 0.6$; the actual curves in this region are as shown in Figure 1.

Often an autocorrelation function could not be fitted because of some deficiency in the data. If the deficiency was only a lack of data about the first negative loop, a partial analysis was carried out. A parabola was fitted by least squares to all the points up to and including the first point whose value of ρ was ≤ -0.08 . Then the value of τ where this parabola intersects the $\rho = -0.08$ line (the dashed line below and parallel to the τ axis inserted for ease in reading values) was used to obtain T_e . This parabola, when used, is plotted through the data points as a dotted curve, and the point of intersection is indicated by a circle (e.g., at heights 395 and 445 km in Figure 3). When even this partial analysis could not be completed, only the data bars are plotted, as for the first 13 heights (145-265 km) in Figure 3.

In the first row of Figure 3 some of the measured ρ 's are plotted using other symbols. The x's seen throughout the data in the first eight heights (145-215 km) indicate points that have been rejected because the space-time geometry of the JRO configuration implies a high probability of contamination by strong echoes from the equatorial electrojet or by a succeeding pulse. The up and down arrow symbols ($\uparrow \downarrow$) are used for data points that exceed the -0.4 to +0.6 plotter range for ρ , which was imposed to avoid confusion with other data.

The ρ values as used in this analysis and presented on the microfilm have been adjusted in a manner that must be discussed in some detail. When the preliminary T_e and T_i values were plotted, it was noticed that at night the ratio T_e/T_i was usually less than unity by a significant factor which was statistically independent of height. If this were true, it would require a large heat input to the ion gas with a special height dependence. Since in the nighttime equatorial F2 layer there is not known to be any large heat input with the required height dependence, we have assumed that this effect was caused by systematic errors.

In order to study these systematic errors, we first restricted our attention to data taken between 1900 and 0300 hours (75° W time) since earlier evening data and later morning data often seemed to be in a transition state between day and nighttime conditions. Heights between 300 and 400 km only were considered in order to reduce the possibility of composition effects. The median T_e/T_i , denoted $(T_e/T_i)_{med}$, within this height range was found for each integration, and the results were plotted versus time.

Examination of this graph showed periods when the T_e/T_i bias seemed to be constant except for statistical fluctuations. For each of these periods the factor RF (ρ Factor) by which the ρ values should be divided so that the median of the $(T_e/T_i)_{med}$'s would become 1.0 was calculated. A graph of the RF values versus date is shown in Figure 4.

This normalization method changes T_i much more than T_e , because ρ_{min} is changed inversely as RF but τ at $\rho = -0.08$ is changed very little. Indeed, the changes in T_e are usually within the estimated standard deviation of T_e , although T_e always decreases for $RF \geq 1$, and vice versa.

This method obviously provides only a first-order correction of T_e/T_i . During some periods there was no nighttime data suitable for a determination of RF. For this reason the last two months of 1966 data and the first two and one-half months of 1967 data had to be assigned the mean RF value (namely, 0.9) of all of the other data. The periods when the nighttime values of T_e/T_i were obtained are indicated at the bottom of Figure 4. Furthermore, the statistical fluctuations about $(T_e/T_i)_{med}$ over the periods for which each RF was determined are not negligible. Therefore, in some studies the user may wish to renormalize T_e/T_i according to a physical model of temperature structure. This can be done approximately by changing the RF for each profile or set of profiles, with the aid of Equation (1). Unless the change in RF is large, the changes in T_e are negligible, that is, all of the changes in T_e/T_i can be attributed to changes in T_i .

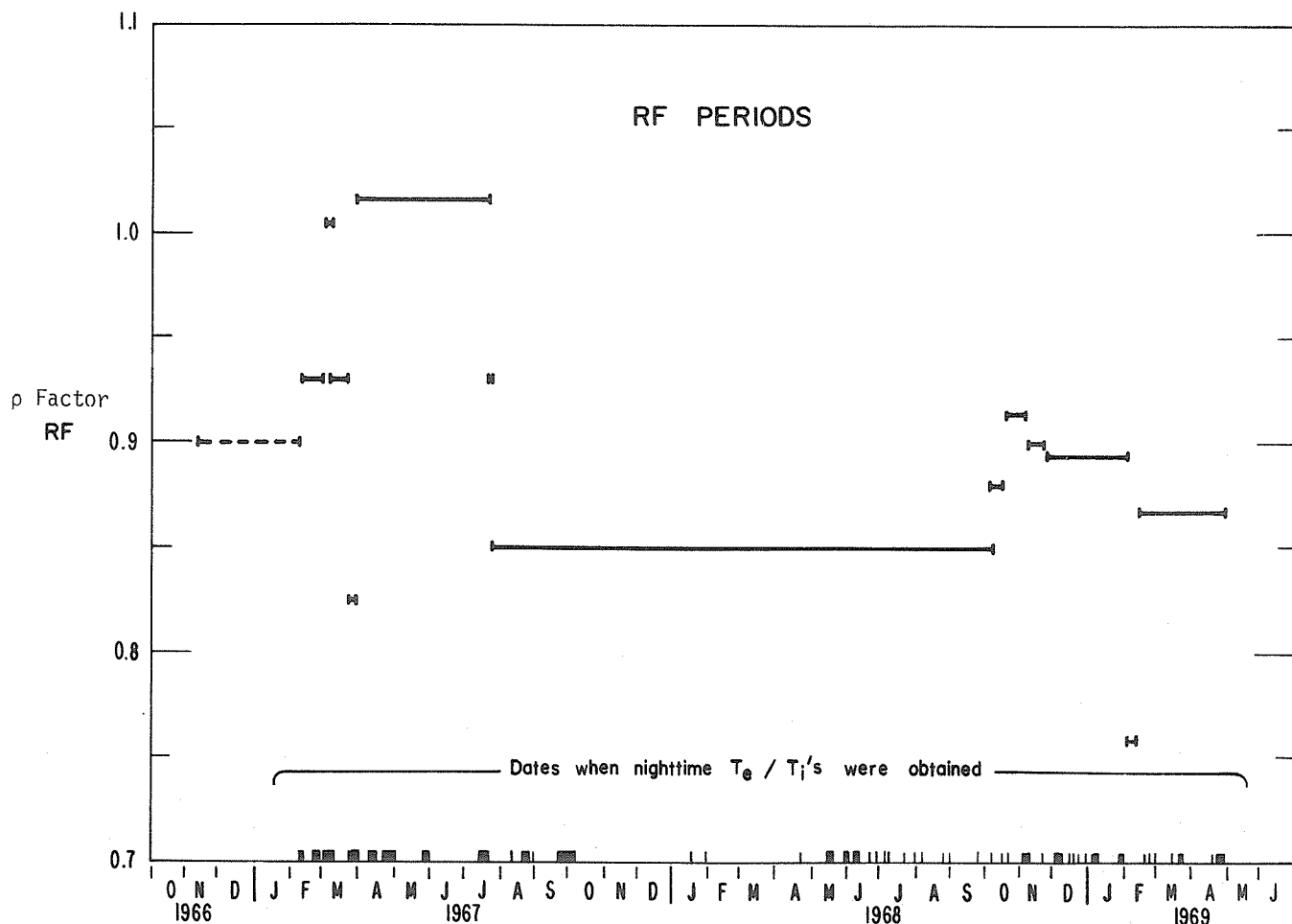


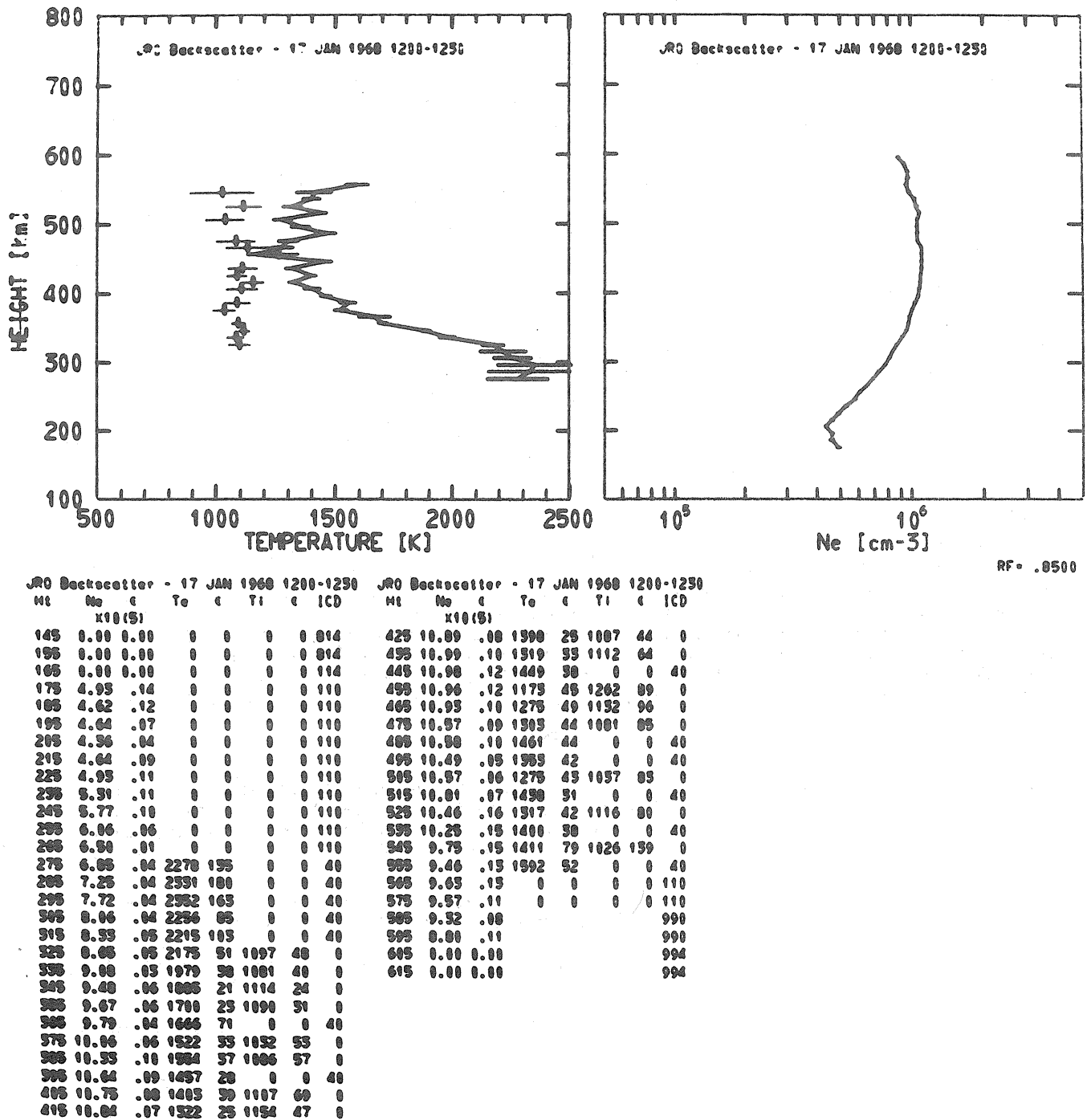
Fig. 4. The RF divisor values applied to the ρ 's plotted over the periods of their use. The periods when nighttime T_e/T_i values were obtained and used in the determination of RF are indicated at the bottom of the graph. The mean RF value of 0.9 was used until 11 February 1967. The process used to determine RF is described in the text.

Format of Archived Data

An example of the graphical and tabular data as it appears on the microfilm is shown in Figure 5. The electron concentration data is presented in the upper right graph and in the second column of the table, labeled Ne. The calculated standard deviation of Ne resulting from the differentiation of the Faraday angle is shown as a horizontal bar on the graph (too small to be visible in this example) and in the first ϵ column, just to the right of the Ne column. It must be emphasized that ϵ is only the standard deviation due to statistical errors, and that the accuracy of Ne can also be affected by systematic errors. In general, rapid variations of Ne with either height or time must be treated with caution. For example, in Figure 5 the rapid increase of Ne with decreasing altitude below 205 km is a systematic error caused by strong scatter from field-aligned irregularities in the E region. At night, when spread-F irregularities are also often present, this kind of error can extend to greater heights. Certain other systematic errors are caused by imperfections in the detectors, particularly at the greatest altitudes where the signal-to-noise ratio is small. Such imperfections probably cause the weak sinusoidal variation in Ne from about 450 to 600 km in Figure 5.

The T_e and T_i data are presented in the upper left graph and in columns four and six, respectively, in the table. On the graph, the T_e points are connected by a solid line and the T_i points are denoted by circles. The estimated standard deviations resulting from statistical errors in the measured values of ρ are shown as horizontal bars on the graph and are listed in the respective ϵ columns (i.e., columns five and seven) in the table. The column labeled ICD is a diagnostic code used in the analysis program and can be ignored.

The original magnetic tape archived in the data center is a 9-track, 1600 phase encoded binary tape. A detailed description of the format is available upon request from WDC-A, and other types of tapes (such as 7-track binary) may be made available. The tape contains all the Ne, T_e , and T_i results, but does not contain the autocorrelation function data from which they were obtained.



RF = .0500

Fig. 5. Te and Ti are plotted against height in the graph at the upper left. Te is plotted using connected error bars (\pm one standard deviation) while Ti is plotted using circles intersected with error bars. Ne is plotted against the same height scale in the graph at the upper right, using connected error bars. The table in the lower part of the Figure presents the same information in tabular form. Zero values indicate that no results were obtained. The ϵ column gives the estimated standard deviation of the parameter in the column just to the left, in the same units. The "ICD" column contains a computer diagnostic code and may be ignored.

Catalog

Values of Ne, Te, and Ti (normalized as discussed above) along with their associated dates, times and heights and estimated errors are archived at WDC-A on magnetic tape and microfilm. Table 1 is a catalog of the observing periods during which measurements were made. Within each observing period there are one or more radar integrations, each of which is an average of several thousand samples received over 5 to 10 minutes during sunrise, 15 to 30 minutes in the daytime, and 15 to 45 minutes at night. The integrations are usually contiguous in time, but occasionally there are gaps caused by operational difficulties. There are a few time discontinuities on the tape, resulting from integrations being entered on the magnetic tape in non-chronological order. The times written on the magnetic tape and microfilm are correct. Approximately 1600 hours of radar data were analyzed.

Reliability of the Data

The reliability of the present data for any particular aeronomical study can be determined only by careful consideration of the statistical and systematic errors described in this report. The Ne data should be adequate for any study consistent with the statistical errors. On the other hand, the limitation to regions of nearly pure O^+ and the normalization of Te/Ti are especially important to the accuracy of the temperature measurements. Judgment of the adequacy of the data for any particular study or of the correctness of a renormalization procedure must, of course, be the responsibility of the user.

Acknowledgements

We thank A. S. Oldfather and R. H. Winkler for their very significant contributions in the preparation of these data for analysis and we acknowledge the long term efforts of R. Cohen, D. T. Farley, Jr., J. L. Green, D. L. Sterling and the staff of the Jicamarca Radar Observatory in making these measurements. The Jicamarca Radar Observatory was a joint operation of the Instituto Geofísico del Peru and the Environmental Sciences Services Administration (ESSA), a predecessor of NOAA. The research at Jicamarca was partially supported by the National Aeronautical and Space Administration (NASA) Fund Transfer R-06-012-008. The National Center for Atmospheric Research (NCAR), which is sponsored by the National Science Foundation, granted the computer time used for the analyses.

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TABLE 1

CATALOG OF NE, TE AND TI MEASUREMENTS AT JICAMARCA RADAR OBSERVATORY

TIME INTERVAL COVERED						TIME INTERVAL COVERED					
YEAR	DATE	75WMT	DATE	75WMT	NO. OF PROFILES	YEAR	DATE	75WMT	DATE	75WMT	NO. OF PROFILES
1966	NOV 11,	0422	TO NOV 11,	1915	50	1967	APR 14,	2145	TO APR 15,	0245	16
	NOV 12,	0315	TO NOV 12,	1215	55		APR 16,	1035	TO APR 16,	1220	8
	NOV 13,	0530	TO NOV 13,	1145	29		APR 17,	1200	TO APR 17,	1330	4
	NOV 14,	1640	TO NOV 16,	0605	121		APR 20,	1345	TO APR 20,	1500	3
	DEC 1,	1210	TO DEC 2,	0430	45		APR 21,	1330	TO APR 21,	2100	15
	DEC 2,	0835	TO DEC 2,	1130	7		APR 25,	0100	TO APR 25,	0415	7
	DEC 7,	1506	TO DEC 7,	1635	2		APR 25,	1215	TO APR 25,	1500	6
	DEC 9,	1353	TO DEC 9,	1619	3		APR 26,	1845	TO APR 26,	2130	6
	DEC 12,	0912	TO DEC 12,	1240	5		APR 27,	0130	TO APR 27,	0330	5
	DEC 13,	1715	TO DEC 14,	0930	50		APR 27,	1210	TO APR 28,	0300	25
	DEC 14,	1500	TO DEC 14,	1600	3		APR 30,	1900	TO APR 30,	2330	10
	DEC 15,	1830	TO DEC 15,	2300	7		MAY 1,	1900	TO MAY 1,	2230	14
	DEC 16,	1615	TO DEC 17,	0300	20		MAY 2,	1020	TO MAY 2,	1605	13
	DEC 19,	1745	TO DEC 20,	0315	20		MAY 3,	1130	TO MAY 3,	1315	4
	DEC 20,	1645	TO DEC 20,	1845	5		MAY 3,	1530	TO MAY 3,	1615	3
	DEC 21,	1625	TO DEC 22,	0245	22		MAY 12,	0000	TO MAY 12,	0530	19
	DEC 22,	1430	TO DEC 22,	1830	10		MAY 12,	1115	TO MAY 12,	1505	33
1967	JAN 2,	1400	TO JAN 2,	1800	9		MAY 16,	1600	TO MAY 16,	1615	2
	JAN 3,	1036	TO JAN 3,	1530	6		MAY 24,	1150	TO MAY 25,	0330	35
	JAN 3,	1900	TO JAN 4,	0330	10		MAY 25,	1045	TO MAY 25,	2100	29
	JAN 4,	1015	TO JAN 5,	0215	20		MAY 26,	0030	TO MAY 26,	0200	4
	JAN 5,	1020	TO JAN 6,	0145	19		MAY 26,	0920	TO MAY 26,	1715	13
	JAN 6,	1410	TO JAN 6,	1845	8		MAY 27,	1645	TO MAY 27,	2330	15
	JAN 10,	1036	TO JAN 10,	2000	145		MAY 29,	1210	TO MAY 29,	1800	109
	JAN 11,	1555	TO JAN 12,	0245	17		MAY 30,	1200	TO MAY 30,	1800	100
	JAN 12,	1745	TO JAN 13,	0245	14		MAY 31,	0118	TO MAY 31,	0130	2
	JAN 13,	1104	TO JAN 14,	0145	25		MAY 31,	0415	TO MAY 31,	0445	3
	JAN 14,	0922	TO JAN 15,	0245	31		MAY 31,	0815	TO MAY 31,	0930	3
	JAN 15,	0830	TO JAN 16,	0245	34		MAY 31,	1645	TO MAY 31,	1845	6
	JAN 19,	1030	TO JAN 20,	0300	29		MAY 31,	2230	TO JUN 1,	0345	7
	JAN 20,	1425	TO JAN 20,	1430	2		JUN 1,	1035	TO JUN 1,	1035	1
	JAN 20,	1800	TO JAN 21,	0300	14		JUN 1,	1345	TO JUN 1,	1400	2
	JAN 22,	1325	TO JAN 22,	2025	84		JUN 1,	1645	TO JUN 1,	1700	2
	JAN 23,	1520	TO JAN 24,	0300	20		JUN 2,	1100	TO JUN 2,	1100	1
	JAN 24,	1022	TO JAN 24,	2200	23		JUN 2,	1515	TO JUN 2,	1600	2
	JAN 25,	0300	TO JAN 25,	0300	1		JUN 8,	1122	TO JUN 8,	1225	2
	JAN 31,	1903	TO JAN 31,	1926	3		JUN 8,	1545	TO JUN 8,	1545	1
	JAN 31,	2200	TO FEB 1,	0840	22		JUN 9,	1545	TO JUN 9,	1545	1
	FEB 2,	0300	TO FEB 2,	1045	27		JUN 12,	0910	TO JUN 12,	0945	2
	FEB 7,	1330	TO FEB 8,	0400	24		JUN 16,	1500	TO JUN 16,	1808	6
	FEB 8,	0930	TO FEB 9,	0500	30		JUN 22,	1530	TO JUN 22,	2315	17
	FEB 9,	0845	TO FEB 9,	1100	8		JUN 24,	1315	TO JUN 24,	1415	3
	FEB 10,	1000	TO FEB 11,	0500	30		JUN 24,	1715	TO JUN 24,	1945	9
	FEB 16,	1015	TO FEB 17,	0245	29		JUN 27,	1145	TO JUN 27,	1215	3
	FEB 17,	0845	TO FEB 18,	0530	44		JUL 6,	1100	TO JUL 6,	1600	9
	FEB 21,	1545	TO FEB 21,	1700	4		JUL 7,	0950	TO JUL 7,	1550	13
	FEB 22,	1130	TO FEB 22,	1505	9		JUL 11,	2300	TO JUL 11,	2300	1
	FEB 23,	0905	TO FEB 23,	0950	4		JUL 12,	0200	TO JUL 12,	0400	2
	FEB 23,	1200	TO FEB 24,	0330	31		JUL 13,	0405	TO JUL 13,	0405	1
	FEB 24,	0830	TO FEB 24,	1315	11		JUL 13,	0725	TO JUL 13,	0725	1
	FEB 24,	1600	TO FEB 25,	0200	16		JUL 13,	1550	TO JUL 13,	1550	1
	FEB 27,	1200	TO FEB 27,	1230	2		JUL 17,	1400	TO JUL 20,	0750	149
	FEB 28,	1130	TO FEB 28,	1230	2		JUL 21,	0400	TO JUL 21,	0400	1
	MAR 1,	0915	TO MAR 1,	1757	120		JUL 21,	0800	TO JUL 21,	0800	1
	MAR 3,	0900	TO MAR 3,	0900	1		JUL 24,	2305	TO JUL 24,	2305	1
	MAR 3,	1145	TO MAR 3,	1330	4		JUL 25,	2000	TO JUL 26,	0800	33
	MAR 3,	1600	TO MAR 4,	0345	25		JUL 31,	0920	TO JUL 31,	0940	2
	MAR 6,	1215	TO MAR 6,	1230	2		JUL 31,	1630	TO JUL 31,	1730	2
	MAR 7,	0930	TO MAR 7,	1230	8		AUG 1,	2230	TO AUG 2,	0200	7
	MAR 22,	1400	TO MAR 22,	1700	6		AUG 2,	0500	TO AUG 2,	0915	12
	MAR 22,	2030	TO MAR 22,	2030	1		AUG 2,	1615	TO AUG 3,	1200	40
	MAR 23,	1100	TO MAR 23,	2230	13		AUG 7,	2040	TO AUG 7,	2040	1
	MAR 24,	0100	TO MAR 24,	0100	1		AUG 9,	2330	TO AUG 9,	2330	1
	MAR 27,	1045	TO MAR 27,	1700	48		AUG 10,	1500	TO AUG 11,	0115	19
	MAR 29,	1400	TO MAR 29,	1400	1		AUG 11,	0935	TO AUG 12,	0750	52
	MAR 29,	1630	TO MAR 29,	1800	2		AUG 14,	1300	TO AUG 17,	1030	157
	MAR 30,	1530	TO MAR 31,	0330	24		AUG 17,	2311	TO AUG 17,	2311	1
	MAR 31,	1145	TO MAR 31,	1230	3		AUG 23,	1145	TO AUG 23,	1515	5
	APR 5,	1000	TO APR 5,	1730	15		AUG 28,	1745	TO AUG 29,	1200	48
	APR 12,	1430	TO APR 13,	0330	27		AUG 29,	1510	TO AUG 31,	0900	104
	APR 13,	1030	TO APR 14,	0115	33		SEP 1,	1450	TO SEP 1,	1940	59
	APR 14,	1115	TO APR 14,	1445	8		SEP 5,	1845	TO SEP 6,	0100	16

TABLE 1 (CONT'D)

CATALOG OF NE, TE AND TI MEASUREMENTS AT JICAMARCA RADAR OBSERVATORY

TIME INTERVAL COVERED						TIME INTERVAL COVERED					
YEAR	DATE	75WMT	DATE	75WMT	NO. OF PROFILES	YEAR	DATE	75WMT	DATE	75WMT	NO. OF PROFILES
1967	SEP 7,	1815	TO SEP 8,	0100	18	1968	JUL 23,	1615	TO JUL 24,	0745	32
	SEP 8,	1428	TO SEP 8,	1645	3		JUL 24,	1145	TO JUL 24,	2345	23
	SEP 12,	1700	TO SEP 13,	0930	36		JUL 30,	0800	TO JUL 31,	0730	52
	OCT 10,	1228	TO OCT 12,	0800	105		AUG 6,	1630	TO AUG 7,	1600	56
	OCT 12,	1200	TO OCT 12,	1230	2		AUG 8,	1415	TO AUG 8,	1430	2
	OCT 24,	1600	TO OCT 24,	2345	16		AUG 9,	1200	TO AUG 9,	1230	3
	OCT 26,	1645	TO OCT 26,	2345	13		AUG 26,	1650	TO AUG 26,	2300	16
	OCT 28,	1700	TO OCT 28,	2300	14		AUG 27,	2145	TO AUG 28,	1530	42
	OCT 30,	1555	TO OCT 30,	1555	1		SEP 3,	1620	TO SEP 4,	0830	36
	NOV 3,	0515	TO NOV 3,	0830	7		SEP 4,	1430	TO SEP 5,	1500	58
	NOV 16,	1115	TO NOV 16,	1400	5		SEP 13,	1200	TO SEP 13,	1315	3
	NOV 16,	1740	TO NOV 16,	2300	10		SEP 23,	1330	TO SEP 23,	1530	4
	NOV 17,	0500	TO NOV 17,	0715	4		SEP 23,	1845	TO SEP 24,	0900	33
	NOV 17,	1245	TO NOV 17,	1245	1		OCT 7,	1330	TO OCT 7,	1745	9
	NOV 20,	1535	TO NOV 20,	1535	1		OCT 8,	1132	TO OCT 9,	1607	57
	NOV 23,	0850	TO NOV 25,	0800	180		OCT 15,	1605	TO OCT 16,	0200	30
	NOV 28,	1815	TO NOV 29,	0130	19		OCT 16,	1515	TO OCT 17,	0200	34
	NOV 29,	0830	TO NOV 29,	0918	2		OCT 18,	1240	TO OCT 18,	1610	21
	NOV 29,	1215	TO NOV 29,	1245	2		OCT 21,	1100	TO OCT 21,	1435	19
	NOV 29,	1745	TO NOV 29,	2300	20		OCT 21,	1720	TO OCT 22,	0800	47
	DEC 12,	1140	TO DEC 12,	1230	3		OCT 22,	1130	TO OCT 22,	1200	2
	DEC 12,	1750	TO DEC 13,	0330	23		OCT 22,	1630	TO OCT 23,	0345	36
	DEC 13,	1845	TO DEC 13,	1945	5		OCT 31,	1315	TO NOV 1,	1200	58
	DEC 18,	1300	TO DEC 19,	2300	84		NOV 1,	1405	TO NOV 1,	2345	22
	DEC 21,	1200	TO DEC 21,	1300	3		NOV 2,	0200	TO NOV 5,	0600	170
	DEC 26,	1200	TO DEC 26,	1300	3		NOV 5,	1015	TO NOV 6,	0500	43
	DEC 27,	1145	TO DEC 27,	1415	6		NOV 6,	1230	TO NOV 6,	1250	2
	DEC 29,	1200	TO DEC 29,	1330	4		NOV 6,	1530	TO NOV 6,	2330	15
1968	JAN 2,	1145	TO JAN 2,	1330	5		NOV 19,	1110	TO NOV 21,	0615	97
	JAN 3,	1200	TO JAN 3,	1230	2		DEC 3,	1245	TO DEC 3,	2045	20
	JAN 4,	1200	TO JAN 4,	1300	3		DEC 3,	2300	TO DEC 4,	1545	40
	JAN 16,	0930	TO JAN 16,	1330	9		DEC 16,	1545	TO DEC 16,	2133	11
	JAN 16,	1625	TO JAN 17,	2345	71		DEC 17,	0930	TO DEC 17,	1100	3
	JAN 28,	2220	TO JAN 28,	2300	3		DEC 17,	1400	TO DEC 18,	0800	45
	FEB 22,	1200	TO FEB 22,	2330	28		DEC 18,	1130	TO DEC 18,	1300	4
	MAR 13,	0855	TO MAR 13,	1330	8		DEC 20,	1515	TO DEC 21,	0000	29
	MAR 13,	1600	TO MAR 14,	0800	53		DEC 23,	1640	TO DEC 24,	0015	23
	APR 2,	1830	TO APR 3,	0830	48		DEC 27,	1530	TO DEC 27,	2345	18
	APR 3,	1400	TO APR 4,	0815	55		DEC 30,	1405	TO DEC 30,	1405	1
	APR 23,	1215	TO APR 24,	0800	40		DEC 30,	1730	TO DEC 30,	2345	15
	APR 24,	1200	TO APR 24,	1300	3	1969	JAN 7,	0700	TO JAN 7,	1003	6
	APR 24,	1645	TO APR 24,	2357	21		JAN 7,	1330	TO JAN 8,	0800	41
	APR 25,	0200	TO APR 25,	0800	17		JAN 28,	1020	TO JAN 28,	1100	3
	APR 26,	0130	TO APR 26,	0800	20		JAN 28,	1345	TO JAN 29,	0800	42
	APR 29,	2215	TO APR 30,	0300	10		JAN 29,	1120	TO JAN 30,	0800	53
	APR 30,	1645	TO APR 30,	2345	20		FEB 17,	1530	TO FEB 19,	1130	105
	MAY 2,	0100	TO MAY 2,	0400	6		FEB 24,	1302	TO FEB 24,	1323	2
	MAY 15,	1745	TO MAY 16,	0745	22		FEB 26,	0830	TO FEB 28,	0330	102
	MAY 16,	1345	TO MAY 16,	1405	3		FEB 28,	1145	TO FEB 28,	1230	3
	MAY 17,	0930	TO MAY 18,	0745	41		MAR 4,	1715	TO MAR 6,	0800	93
	JUN 5,	1230	TO JUN 6,	1430	44		MAR 17,	1645	TO MAR 18,	1100	62
	JUN 11,	0945	TO JUN 12,	0830	44		MAR 18,	1650	TO MAR 19,	0900	62
	JUN 20,	1250	TO JUN 21,	1500	54		MAR 19,	1640	TO MAR 20,	0900	60
	JUN 24,	1030	TO JUN 24,	1200	5		MAR 25,	1145	TO MAR 26,	0830	54
	JUN 24,	1630	TO JUN 24,	2245	13		MAR 27,	1315	TO MAR 27,	1600	4
	JUN 25,	1245	TO JUN 26,	1550	67		APR 3,	0945	TO APR 3,	1100	5
	JUL 1,	1245	TO JUL 1,	1300	2		APR 10,	1805	TO APR 10,	2345	24
	JUL 8,	1645	TO JUL 8,	1705	2		APR 11,	1920	TO APR 11,	2245	14
	JUL 9,	0800	TO JUL 10,	1630	69		APR 12,	2020	TO APR 12,	2230	10
	JUL 11,	1500	TO JUL 11,	1500	1		APR 13,	2040	TO APR 13,	2345	14
	JUL 12,	1440	TO JUL 12,	1440	1		APR 14,	2030	TO APR 14,	2315	12
	JUL 16,	1515	TO JUL 16,	1530	2		APR 15,	1845	TO APR 15,	2345	21
	JUL 17,	1220	TO JUL 17,	1250	2		APR 16,	2045	TO APR 16,	2230	8
	JUL 22,	1300	TO JUL 22,	1450	4		APR 22,	1620	TO APR 24,	1500	115
	JUL 23,	1100	TO JUL 23,	1300	5		APR 28,	1400	TO APR 29,	1130	50

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